

Gridless Sound Field Decomposition Based on Reciprocity Gap Functional in Spherical Harmonic Domain

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Abstract

Sound field decomposition

- Goal is to interpolate and reconstruct sound field inside region including sources (ill-posed problem!)
- Sound field should be decomposed into fundamental solutions of Helmholtz eq., i.e., point sources

Current method: Sparse sound field decomposition^[1]

- Sparse decomposition by discretizing possible source region into grid points
- Discrete set of point-source dictionary causes off-grid problem
- High computational cost and large memory are required for large set of grid points

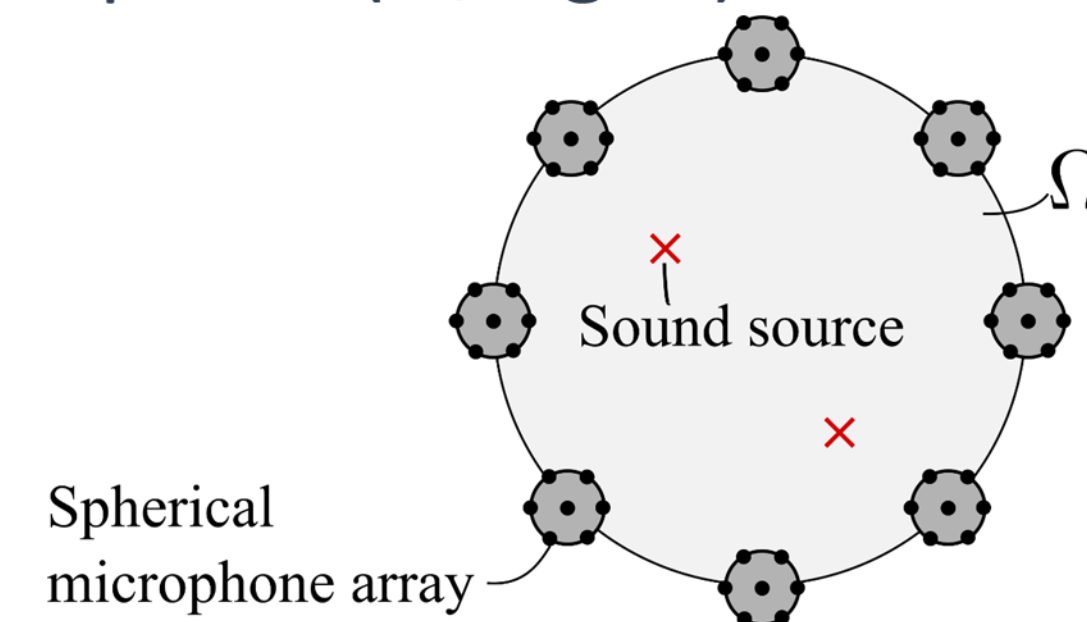
Proposed method

- Reciprocity gap functional (RGF) is applied in spherical harmonic domain
- No more discretization by grids
- Closed-form solution (extremely low computational cost)

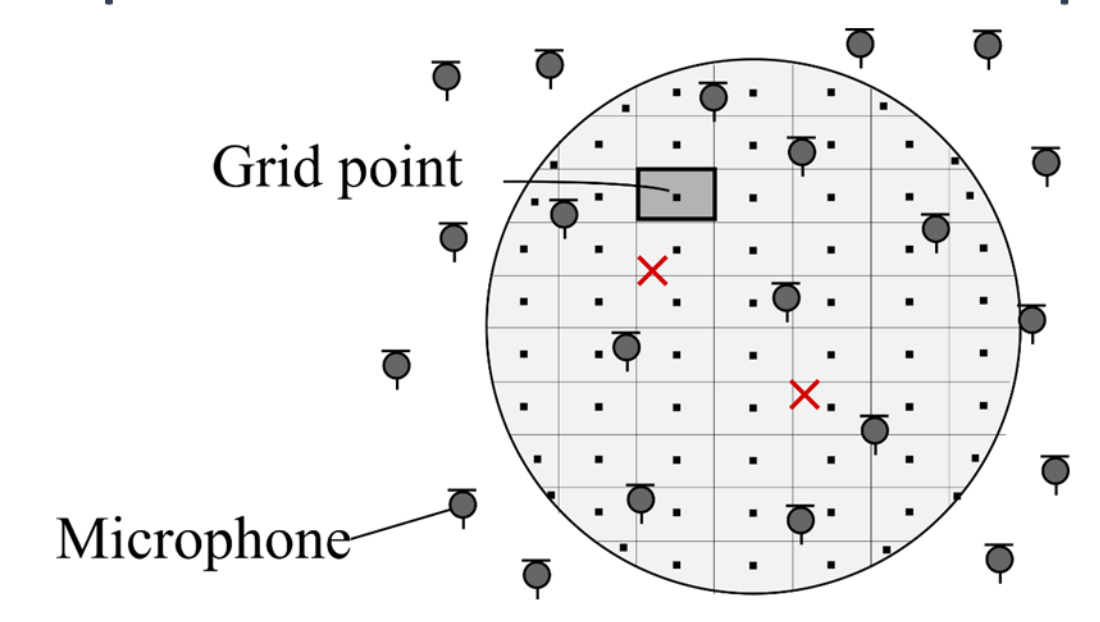
Problem Statement

Sound field decomposition

➤ Proposed (w/o grid)



➤ Sparse sound field decomposition



- Source distribution $Q(\mathbf{r}, k)$ inside Ω is approximated as a linear combination of J point sources

$$Q(\mathbf{r}, k) \approx \sum_{j=1}^J c_j \delta(\mathbf{r} - \mathbf{r}_j) \quad \Rightarrow \quad u(\mathbf{r}, k) \approx \sum_{j=1}^J c_j(k) G(\mathbf{r} | \mathbf{r}_j, k)$$

Spatial convolution of source distribution $Q(\cdot)$ with three-dimensional free-field Green's function $G(\cdot)$

Estimating c_j and \mathbf{r}_j makes it possible to reconstruct $u(\cdot)$

Sound Field Decomposition Based on RGF

Concept of RGF

- Test function $w(\cdot)$ and RGF $R(\cdot)$ for $w(\cdot)$ is defined as

$$w(\cdot) \in \mathcal{H}_k := \{w \in \mathcal{H}^1(\Omega) | (\nabla^2 + k^2)w = 0\}$$

$$R(w) := \int_{\Omega} w(\mathbf{r}) Q(\mathbf{r}) d\mathbf{r}$$

- By applying point source assumption and Green's theorem to $R(w)$, the following equation holds

$$\sum_{j=1}^J c_j w(\mathbf{r}_j) = \int_{\partial\Omega} \left(u(\mathbf{r}) \frac{\partial w(\mathbf{r})}{\partial \mathbf{n}} - w(\mathbf{r}) \frac{\partial u(\mathbf{r})}{\partial \mathbf{n}} \right) dS \quad \dots (*)$$

The parameter of sound sources J , \mathbf{r}_j and c_j can be estimated from pressure and velocity values on $\partial\Omega$.

Localization based on RGF

- $w(\cdot)$ is chosen as $w_n(\mathbf{r}) = (x + iy)^n e^{-ikz}$ [2]
- Calculate s_n (Right-hand side of Eq. (*)) from microphone observations
- Compose Hankel matrices

$$\mathbf{H} = \begin{bmatrix} s_1 & s_2 & \dots & s'_J \\ s_2 & s_3 & \dots & s'_{J+1} \\ \vdots & \vdots & \ddots & \vdots \\ s_{J'} & s_{J'+1} & \dots & s'_{2J'-1} \end{bmatrix}, \quad \mathbf{H}_l = \begin{bmatrix} s_{l+1} & s_{l+2} & \dots & s_{l+J'} \\ s_{l+2} & s_{l+3} & \dots & s_{l+J'+1} \\ \vdots & \vdots & \ddots & \vdots \\ s_{l+J'} & s_{l+J'+1} & \dots & s_{l+2J'-1} \end{bmatrix}$$

Determine # of sources J from the eigenvalue of \mathbf{H}

Estimate the source locations on x - y plane from the eigenvalue of $\mathbf{H}_l \mathbf{H}_l^{-1}$

A large number of microphone pairs need to be arranged uniformly on $\partial\Omega$ to calculate surface integral.

RGF in Spherical Harmonic Domain

RGF in spherical harmonic domain

- Pressure on $\partial\Omega$ in spherical harmonic domain.

$$u(\mathbf{r}) = \sum_{\nu=0}^{\infty} \sum_{\mu=-\nu}^{\nu} \beta_{\nu\mu} h_{\nu}(kr) Y_{\nu\mu}(\theta, \phi)$$

- Spherical harmonic expansion of $w_n(\cdot)$ is analytically obtained

$$w_{\nu\mu}^{(n)} = \frac{1}{j_{\nu}(kr)} \int_0^{\pi} \int_0^{2\pi} w_{\nu\mu}^{(n)} Y_{\nu\mu}^*(\theta, \phi) \sin \theta d\theta d\phi$$

$$= \frac{i^{n-\nu}}{k^n} \sqrt{\frac{4\pi(2\nu+1)(\nu+n)!}{(\nu-n)!}} \delta_{\mu,n}$$

Surface integral in spherical harmonic domain

$$s_n = \int_{\partial\Omega} \left(u(\mathbf{r}) \frac{\partial w(\mathbf{r})}{\partial \mathbf{n}} - w(\mathbf{r}) \frac{\partial u(\mathbf{r})}{\partial \mathbf{n}} \right) dS$$

$$= \frac{i}{k} \sum_{\nu=0}^{\infty} \sum_{\mu=-\nu}^{\nu} (-1)^{\mu+1} w_{\nu-\mu}^{(n)} \beta_{\nu\mu}$$

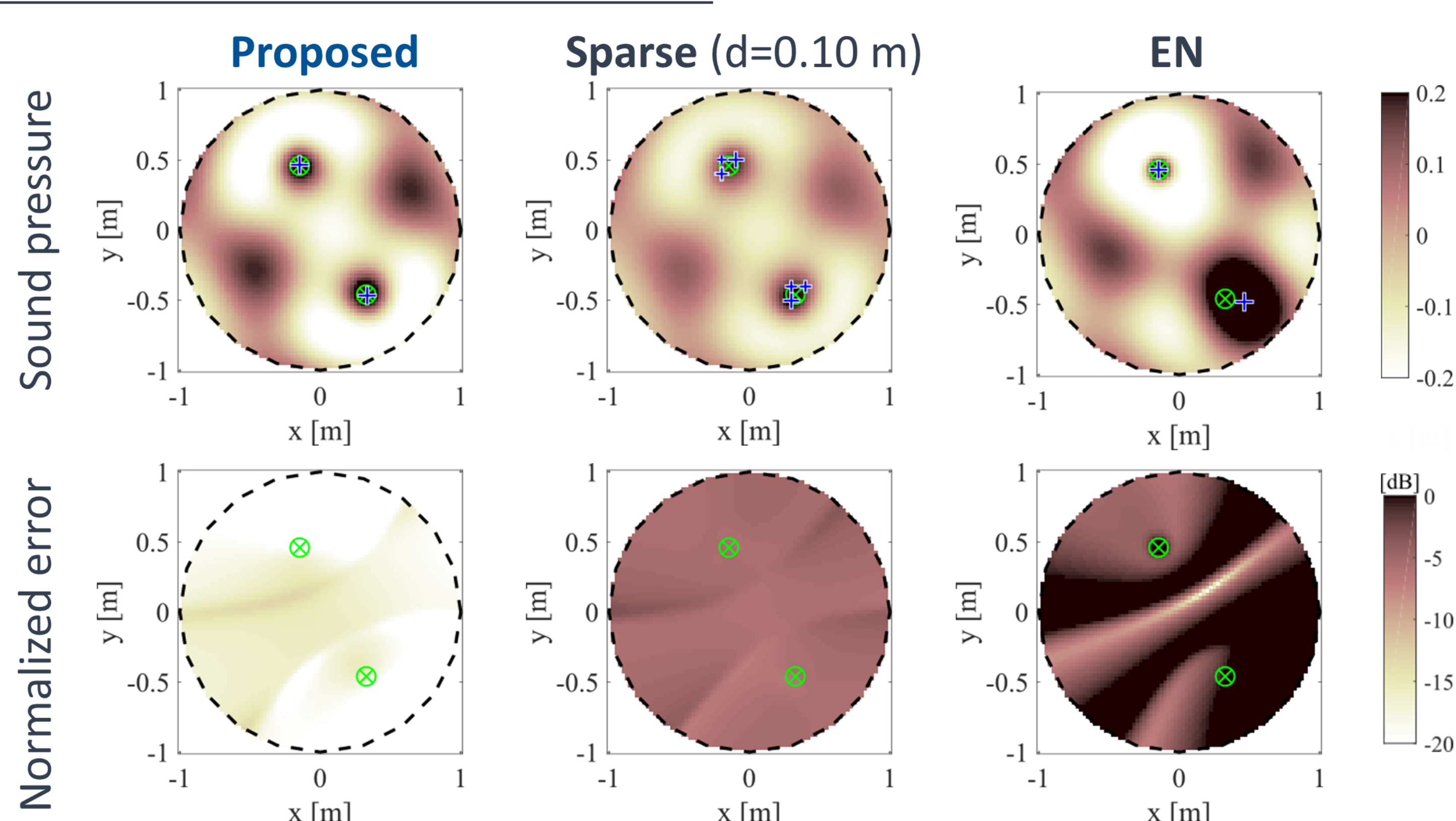
- Two test functions are also used to determine the source location in a 3D space: $w'_n(\mathbf{r}) := w_n(y, z, x)$, $w''_n(\mathbf{r}) := w_n(z, x, y)$
- Analytical spherical harmonic expansion of function $w'_n(\mathbf{r})$, $w''_n(\mathbf{r})$ are obtained from coefficients $w_{\nu\mu}^{(n)}$ and Wigner-D matrix

Experiments

Simulation conditions

- Comparing proposed method (Proposed) with
 - Sparse: sparse sound field decomposition^[1] (interval of grid points d of 0.10, 0.15, and 0.20 m)
 - EN: the direct application of RGF^[2]
- Microphone array (SNR: 40dB)
 - 24 second-order spherical microphone arrays consisting of 9 microphones for Proposed and Sparse
 - 108 microphone pairs for EN

Sound field reconstruction: 400 Hz

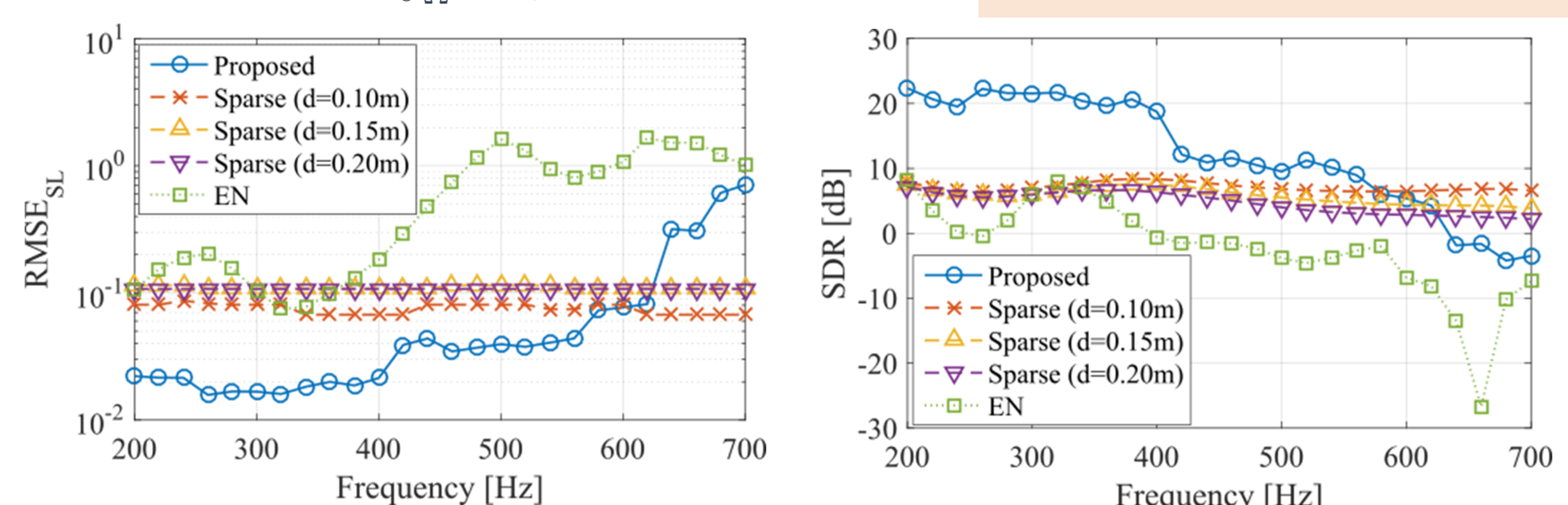


RMSE_{SL} and SDR are plotted against frequency

- Evaluation criteria:
 - Root-mean-square error of source location (RMSE_{SL})
 - Signal-to-distortion ratio (SDR)

$$\text{RMSE}_{\text{SL}} = \sqrt{\frac{1}{J} \sum_{j=1}^J \|\mathbf{r}_{j,\text{true}} - \mathbf{r}_j\|_2^2}$$

$$\text{SDR} = 10 \log_{10} \frac{\int_{\Omega} |u_{\text{P,true}}(\mathbf{r})|^2 d\mathbf{r}}{\int_{\Omega} |u_{\text{P,true}}(\mathbf{r}) - \hat{u}_{\text{P}}(\mathbf{r})|^2 d\mathbf{r}}$$



Computational time (avg.)

- Proposed: 1.97×10^{-4} s
- Sparse
 - d=0.10 m: 1.64×10^2 s
 - d=0.15 m: 1.88×10^1 s
 - d=0.20 m: 5.80 s
- EN: 1.77×10^{-4} s

High source localization and sound field reconstruction accuracies are achieved at low frequencies by using proposed method

[1] S. Koyama, S. Shimauchi, and H. Ohmuro, "Sparse sound field representation in recording and reproduction for reducing spatial aliasing artifacts", in *Proc. IEEE Int. Conf. Acoust., Speech, Signal Process. (ICASSP)*, Florence, May 2014, pp. 4443–4447.
[2] A. El Badia and T. Nara, "An inverse source problem for Helmholtz's equation from the Cauchy data with a single wave number," *Inverse Prob.*, vol. 27, no. 105001, 2011.